

Parity Violating Asymmetry and Dipole Polarizability in ^{208}Pb

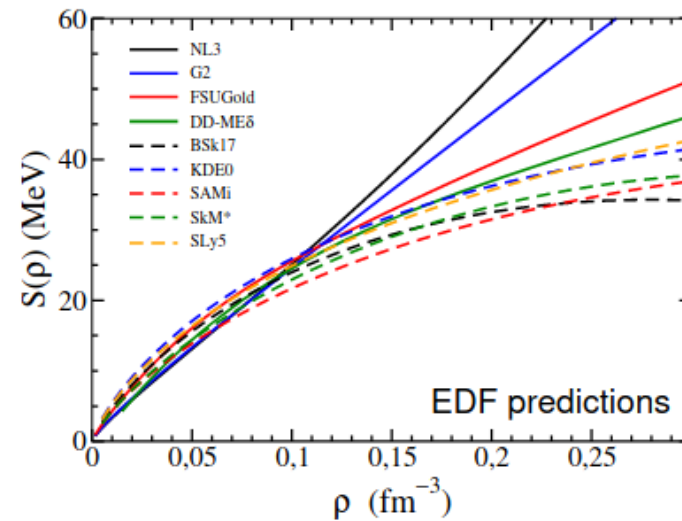
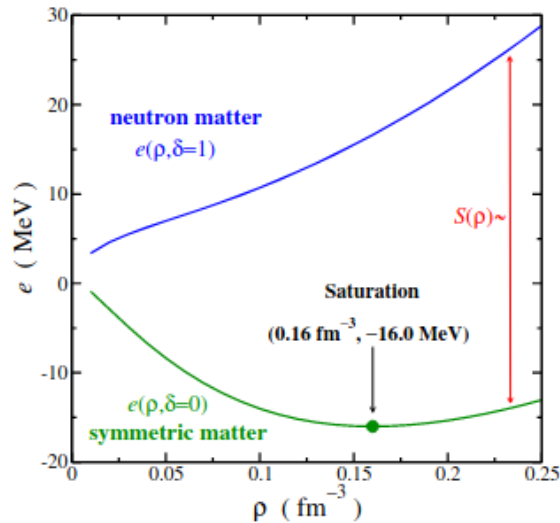
Xavier Roca Maza
Università degli studi di Milano & INFN

*13th International Spring Seminar on Nuclear Physics Perspectives and Challenges in
Nuclear Structure after 70 Years of Shell Model,
Sant'Angelo d'Ischia, May 15-20, 2022*

Table of contents

- **Motivation**
- **Simple models for A_{pv} and α_D :** hints for physical correlations
- **Recent experimental results:** PREx (Jlab) and α_D (RCNP) in ^{208}Pb
- **Previous analysis** based on **correlations** shown within different **EDFs**
- **New analysis** taking into account also **theoretical error** quantification within each **EDF**
- **Conclusions**

Motivation: nuclear EoS



$$e(\rho, \beta) = e(\rho, \beta = 0) + S(\rho)\beta^2 + \mathcal{O}(\beta^4) \quad \text{where} \quad \beta \equiv \frac{\rho_n - \rho_p}{\rho}$$

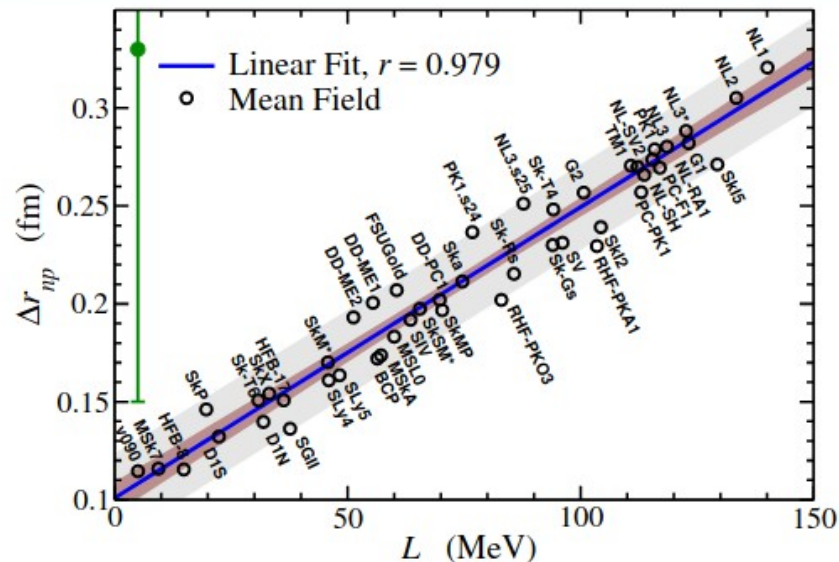
$$J \equiv S(\rho_0); \quad L = 3\rho_0 \left. \frac{\partial S(\rho)}{\partial \rho} \right|_{\rho=\rho_0}$$

-**Isovector properties** not well determined in current **EDFs**
EDFs typically fitted to masses and charge radii

Motivation: neutron skin thickness

But how we can better constraint the isovector channel from observables? (Example)

Neutron skin thickness → is one of the most paradigmatic example of an **isovector sensitive observable**.



But the **neutron skin thickness** is **not** directly **observable**

Physical Review Letters **106**, 252501 (2011)

$$\Delta r_{np} \equiv \langle r_n^2 \rangle^{1/2} - \langle r_p^2 \rangle^{1/2} \sim \frac{1}{12} \frac{N-Z}{A} \frac{R}{J} L ; \quad J \equiv S(\rho_0) \text{ and } L \equiv 3\rho_0 p_0^{\text{neut}}$$

The neutron skin can be probed by the Parity Violating Asymmetry

- Electrons interact by exchanging a γ (couples to p) or a Z_0 boson (couples to n)
- Ultra-relativistic electrons, depending on their helicity ($\pm$), will interact with the nucleus seeing a slightly different potential: Coulomb $\pm$ Weak

- $$A_{pv} \equiv \frac{d\sigma_+/d\Omega - d\sigma_-/d\Omega}{d\sigma_+/d\Omega + d\sigma_-/d\Omega} \sim \frac{\text{Weak}}{\text{Coulomb}}$$

- Input for the calculation are the ρ_p and ρ_n (main uncertainty) and nucleon form factors for the e-m and the weak neutral current.

Electron motion governed by the Dirac equation:

$$[\vec{\alpha} \cdot \vec{p} + \beta m_e + V(r)]\psi = E\psi$$

where $V(r) = V_C(r) + \gamma^5 V_W(r)$

Dirac equation for helicity states ($m_e \approx 0$)

$$[\vec{\alpha} \cdot \vec{p} + (V_C(r) \pm V_W(r))]\psi_{\pm} = E\psi_{\pm}$$

Parity Violating Asymmetry:

definition and simple model (PWBA; $F_W \approx F_n$; and $F_{ch} \approx F_p$)

Within the Plane Wave Born Approximation,

$$A_{pv} = \frac{G_F q^2}{4\pi\alpha\sqrt{2}} \left[4 \sin^2 \theta_W + \frac{F_n(q) - F_p(q)}{F_p(q)} \right]$$

... which depends on $F_n(q) - F_p(q)$. For $q \rightarrow 0$, it is approximately,

$$\begin{aligned} -\frac{q^2}{6} (\langle r_n^2 \rangle - \langle r_p^2 \rangle) &= -\frac{q^2}{6} [\Delta r_{np} (\langle r_n^2 \rangle^{1/2} + \langle r_p^2 \rangle^{1/2})] \\ &= -\frac{q^2}{6} (2\langle r_p^2 \rangle^{1/2} \Delta r_{np} + \Delta r_{np}^2) \end{aligned}$$

variation of A_{pv} at a fixed q dominated by the variation of Δr_{np} . $F_p(q)$ well fixed by experiment

Dipole Polarizability is also a sensitive isovector indicator: definition

The **linear response** or dynamic polarizability of a **nuclear system excited** from its g.s., $|0\rangle$, to an excited state, $|\nu\rangle$, due to the **action of an external isovector oscillating field** (dipolar in our case) of the form $(F e^{i\omega t} + F^\dagger e^{-i\omega t})$:

$$F_{JM} = \sum_i^A r^J Y_{JM}(\hat{r}) \tau_z(i) \quad (\Delta L = 1 \rightarrow \text{Dipole})$$

is proportional to the **static polarizability** for small oscillations

$$\alpha = (8\pi/9) e^2 m_{-1} = (8\pi/9) e^2 \sum_\nu |\langle \nu | F | 0 \rangle|^2 / E \quad \text{where } m_{-1} \text{ is}$$

the inverse energy weighted moment of the **strength function**, defined as, $S(E) = \sum_\nu |\langle \nu | F | 0 \rangle|^2 \delta(E - E_\nu)$

Dipole Polarizability: simple model

electric polarizability measures tendency of the nuclear charge distribution to be distorted ($\alpha \sim \frac{\text{electric dipole moment}}{\text{external electric field}}$)

- ▶ The dielectric theorem establishes that the m_{-1} moment can be computed from the expectation value of the Hamiltonian in the constrained ground state $\mathcal{H}' = \mathcal{H} + \lambda \mathcal{D}$.

Adopting the Droplet Model:

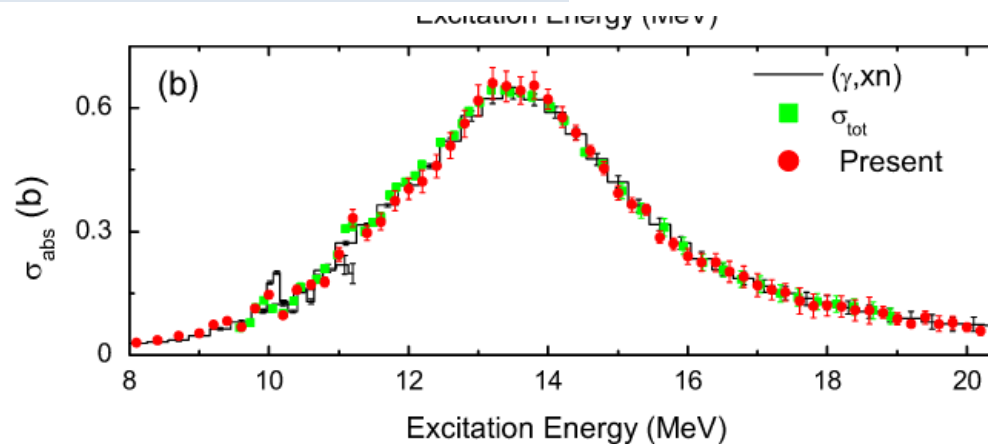
$$m_{-1} \approx \frac{A \langle r^2 \rangle^{1/2}}{48J} \left(1 + \frac{15}{4} \frac{J}{Q} A^{-1/3} \right)$$

within the same model, connection with the neutron skin thickness:

$$\alpha_D \approx \frac{A \langle r^2 \rangle}{12J} \left[1 + \frac{5}{2} \frac{\Delta r_{np} + \sqrt{\frac{3}{5}} \frac{e^2 Z}{70J} - \Delta r_{np}^{\text{surface}}}{\langle r^2 \rangle^{1/2} (I - I_C)} \right]$$

Recent exp. results on A_{pv} and α_D in ^{208}Pb

Tamii et al. PRL107 062502 (2011)



$$\alpha_D = \frac{\hbar c}{2\pi^2 e^2} \int \frac{\sigma_{\text{abs}}}{\omega^2} d\omega,$$

$$\alpha_D(^{208}\text{Pb}) = 19.6 \pm 0.6 \text{ fm}^3$$

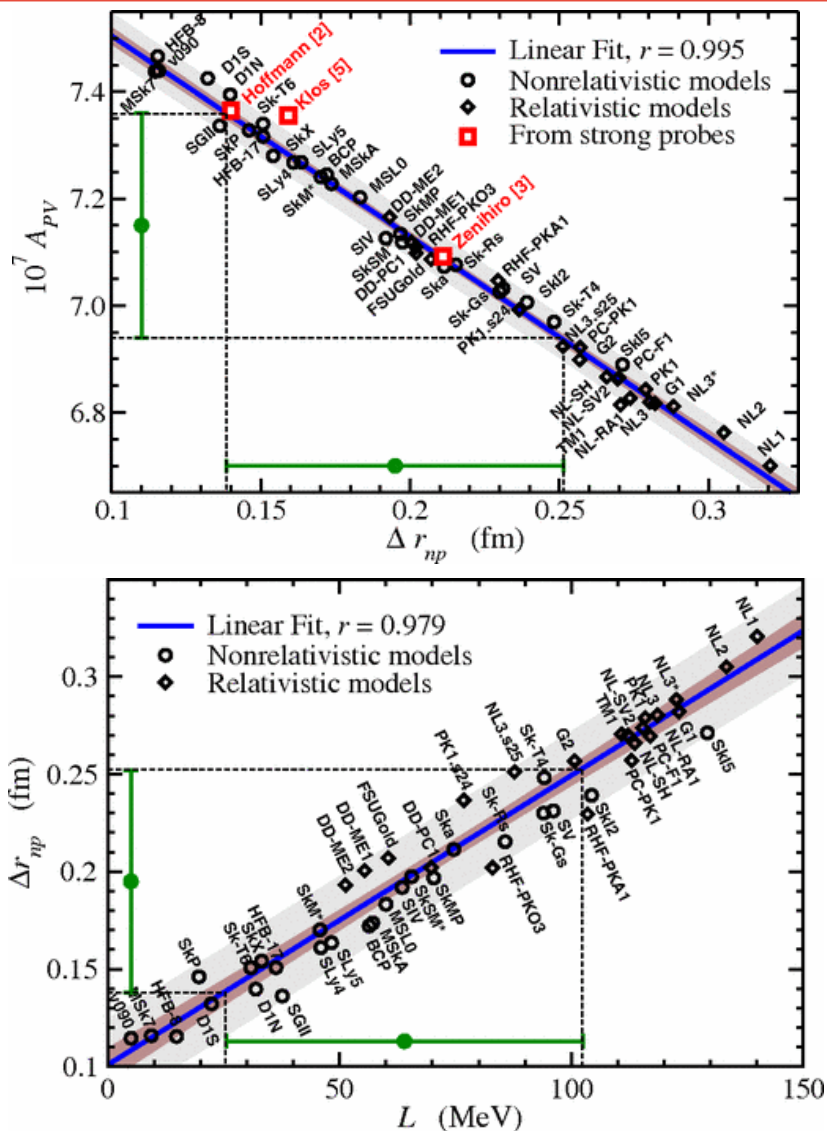
PREx: D. Adhikari et al. (PREX Collaboration) Phys. Rev. Lett. 126, 172502 (2021)

Our final results for $A_{\text{PV}}^{\text{meas}}$ and F_W with the acceptance described by $\epsilon(\theta)$ and $\langle Q^2 \rangle = 0.00616 \text{ GeV}^2$ are

$$A_{\text{PV}}^{\text{meas}} = 550 \pm 16 (\text{stat}) \pm 8 (\text{syst}) \text{ ppb}$$

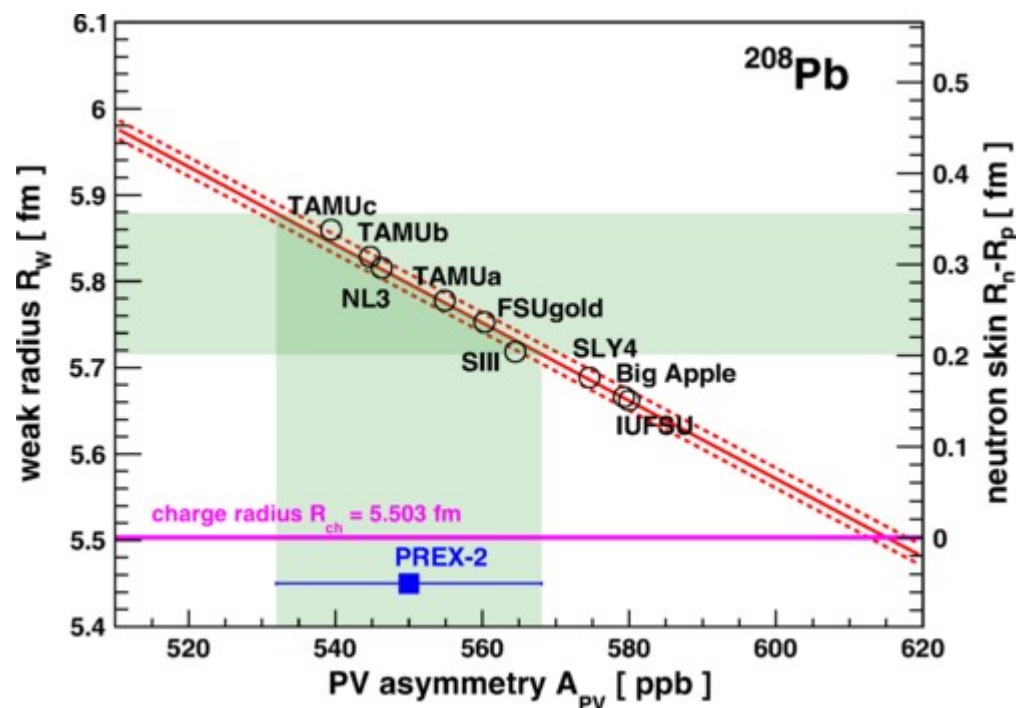
$$F_W(\langle Q^2 \rangle) = 0.368 \pm 0.013 (\text{exp}) \pm 0.001 (\text{theo}),$$

Analysis based on EDFs correlations between A_{pv} , ΔR_{ch} and L



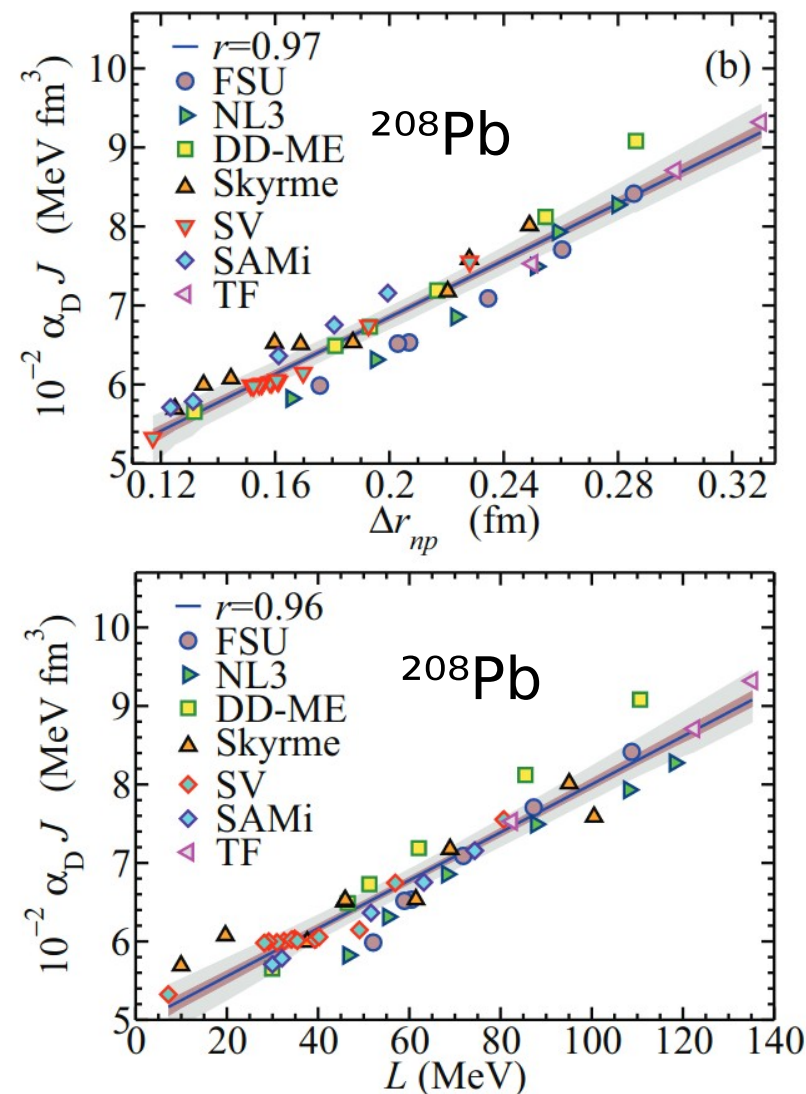
EDFs calibrated to **reproduce** the **binding energy** and **charge radii** of some **selected nuclei**. (In some cases pseudo-observables related to the EoS are also used)

$$R_n - R_p = 0.283 \pm 0.071 \text{ fm.}$$



D. Adhikari et al. Phys. Rev. Lett. 126, 172502 (2021)

Analysis based on EDFs correlations between α_D , ΔR_{ch} and L



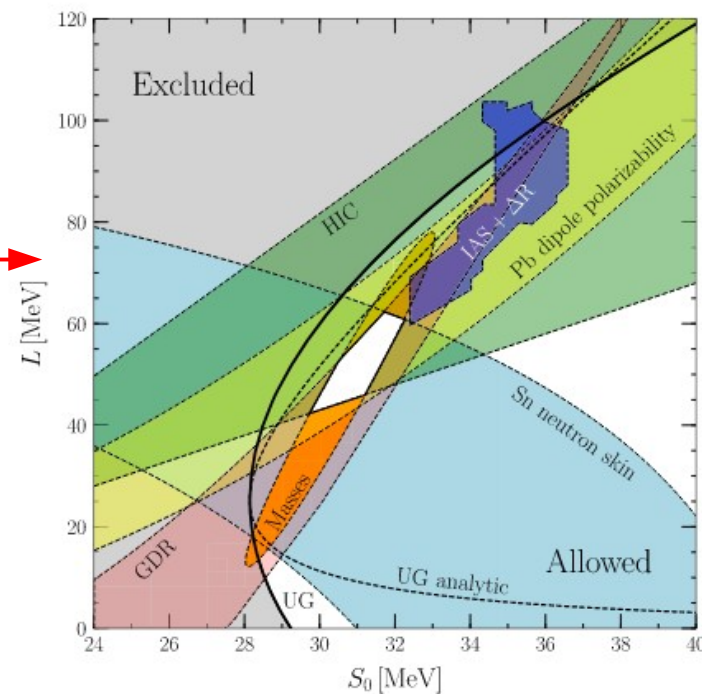
X. Roca-Maza, et al. Phys. Rev. C 88, 024316 (2013)

$$\Delta r_{np} = 0.165 \pm (0.009)_{\text{expt}} \pm (0.013)_{\text{theor}} \pm (0.021)_{\text{est}} \text{ fm.}$$

$$J = (24.5 \pm 0.8) + (0.168 \pm 0.007)L.$$

X. Roca-Maza, et al. Phys. Rev. C 92, 064304 (2015)

Typical plot containing different constraints from different analysis performed in different ways on two non-observable quantities



Ingo Tews et al 2017 ApJ 848 105

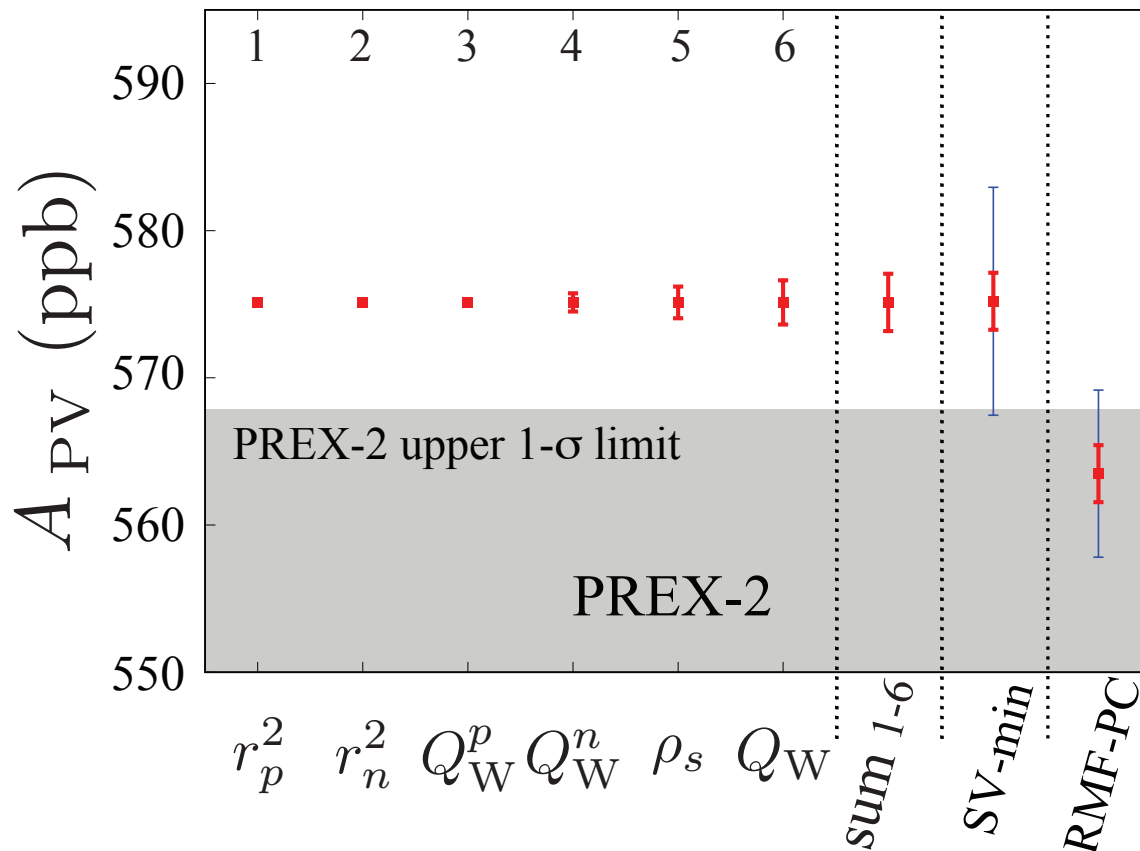
New analysis taking into account model error quantification

How to improve current analysis?

- Include theoretic **statistical errors** and **correlations** within a given **EDF** parametrization
- **Fitting** procedure including experimental data not only on **B** and **R_{ch}** but also on **A_{p,v}** and/or **α_b** (**“informed” EDFs**)
- ...
- **Generalizing** available **EDFs**?
- Other issues related to theory?

PREx: theo. uncertainty budget for A_{pv}

Uncertainties in the determination of the Form Factors is smaller than typical EDFs statistical uncertainties



Parameter	Value
$\langle r_p^2 \rangle$ (fm ²)	0.726 ± 0.019
$\langle r_n^2 \rangle$ (fm ²)	-0.1161 ± 0.0022
μ_p	2.792 847
μ_n	-1.9130
$Q_p^{(W)}$	0.0713 ± 0.0001
$Q_n^{(W)}$	-0.9888 ± 0.0011
ρ_s	-0.24 ± 0.70
κ_s	-0.017 ± 0.004
$Q_{126,82}^{(W)}$	-117.9 ± 0.3

Thin blue bars: statistical model uncertainties (related to neutron and proton densities)

Paul-Gerhard Reinhard, Xavier Roca-Maza, and Witold Nazarewicz *Phys. Rev. Lett.* 127, 232501 (2021)

EDF predictions with error ellipsoids for A_{pv} and α_D in ^{208}Pb

Correlation ellipsoids within each **EDF** show **similar correlations** than the **systematic** study with many **EDFs** $\leftrightarrow$ **have we learnt something?**

FSU $\rightarrow$ Non linear Walecka model (Rel)

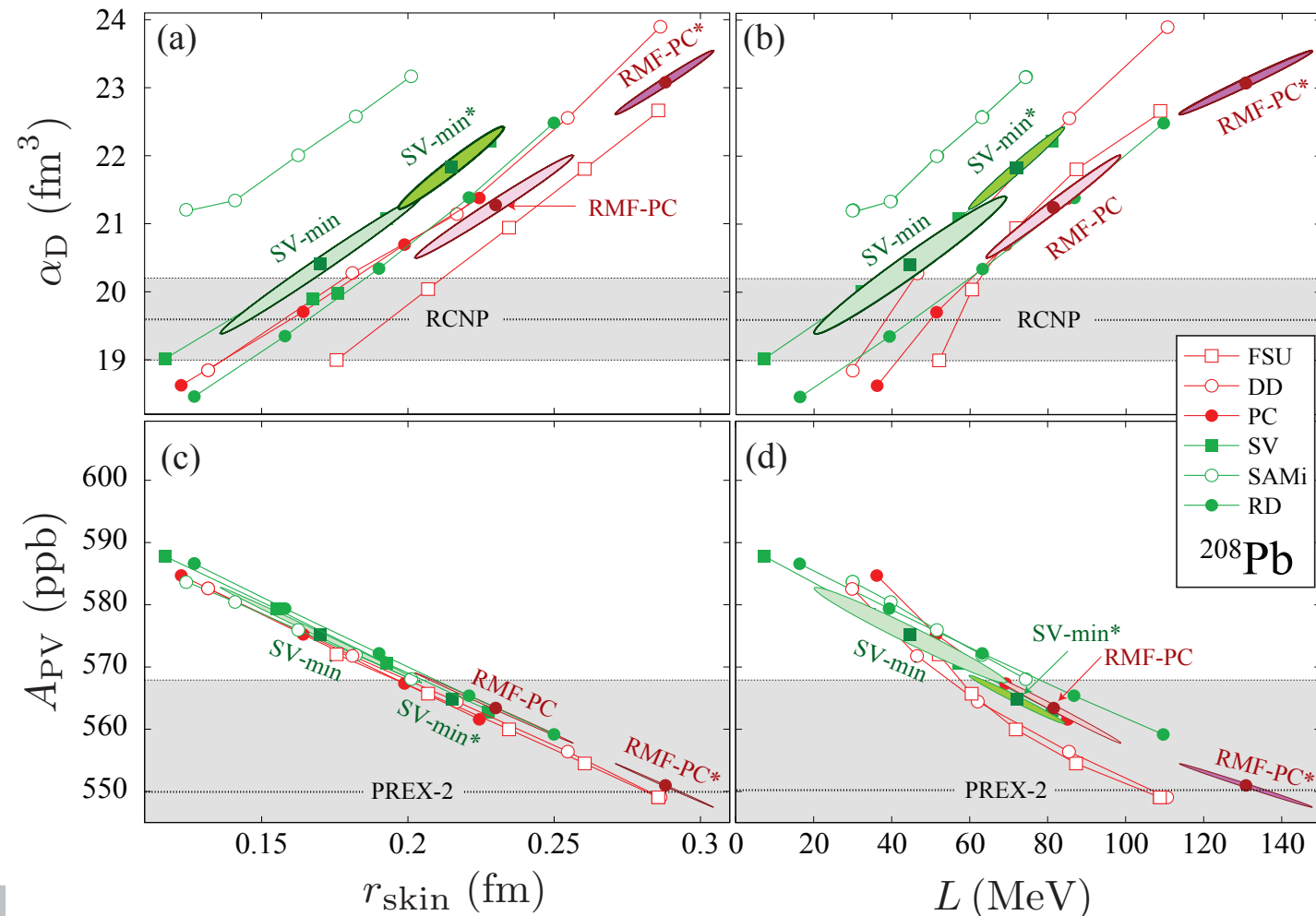
DD $\rightarrow$ Meson-exchange with density dependent couplings (Rel)

PC $\rightarrow$ zero-range density dependent couplings (Rel)

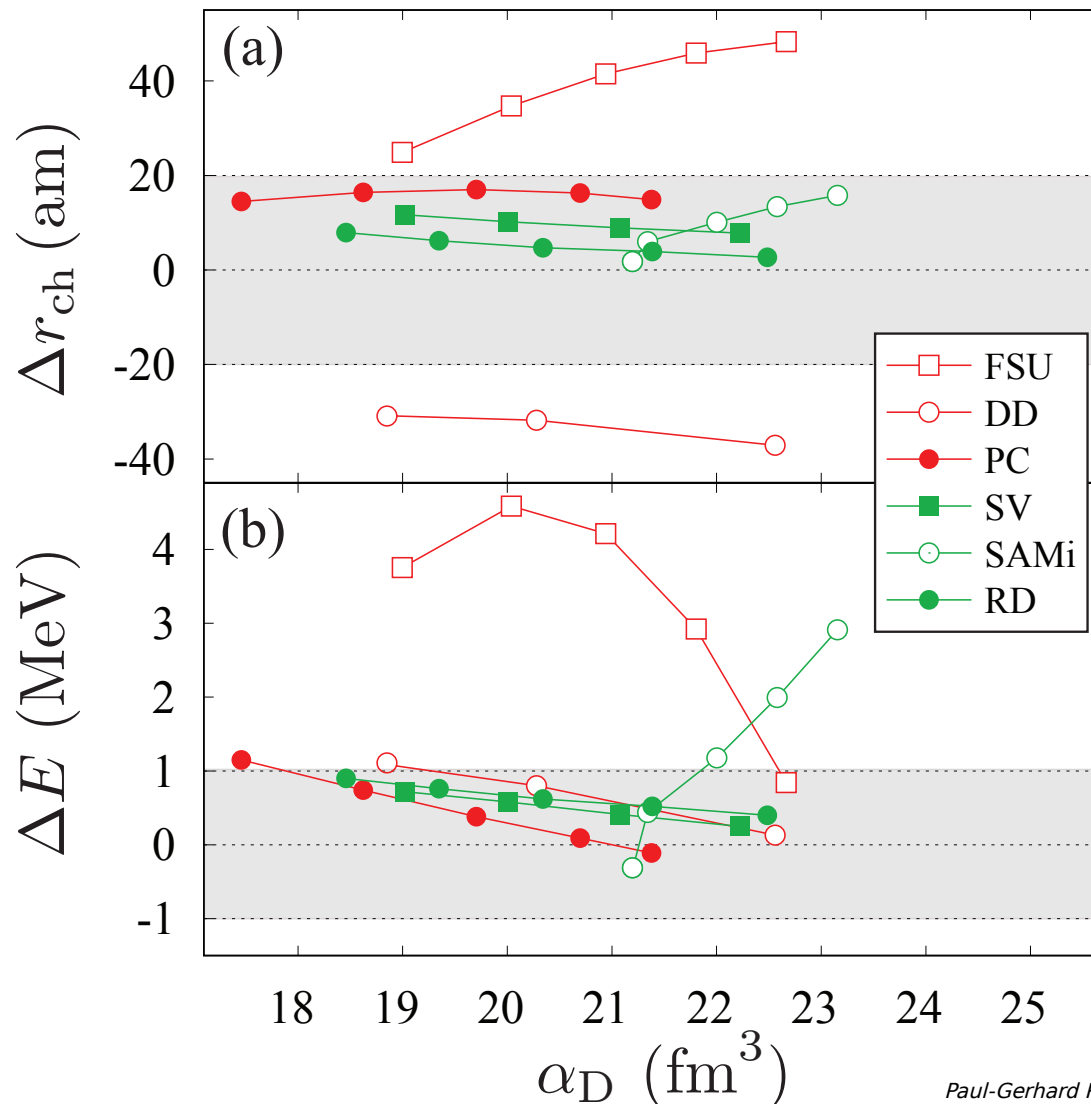
SV&SAMI $\rightarrow$ Skyrme (zero-range, Non-rel)

RD $\rightarrow$ Skyrme with modified density dependence ($\rho^{\text{integer powers}}$)

Models with a "*" contain A_{pv} in the fitting protocol



How these models perform for B and R_{ch} in ^{208}Pb ?

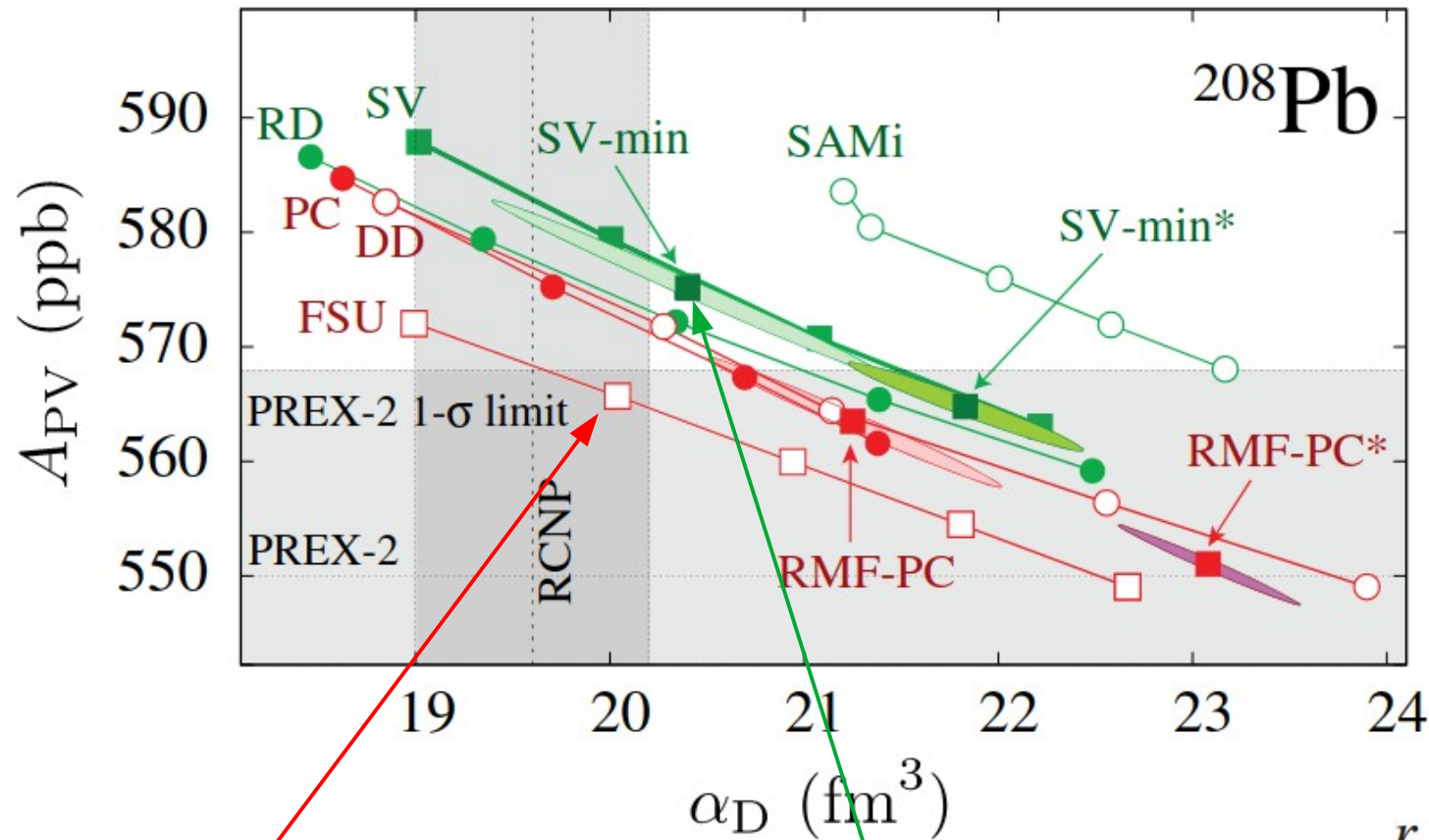


→ The **residuals** of the charge **radius** (a) and binding **energy** (b) of ^{208}Pb for the theoretical models.

→ The **grey bands** around the perfect match indicate the typical performance of EDFs (i.e. **typical r.m.s. deviation** taken over all nuclei where correlation effects are small)

Paul-Gerhard Reinhard, Xavier Roca-Maza, and Witold Nazarewicz Phys. Rev. Lett. 127, 232501 (2021)

A_{pv} versus α_D in well calibrated EDFs



SV-min, suitable model to predict EoS around ρ_0 and Δr_{np}

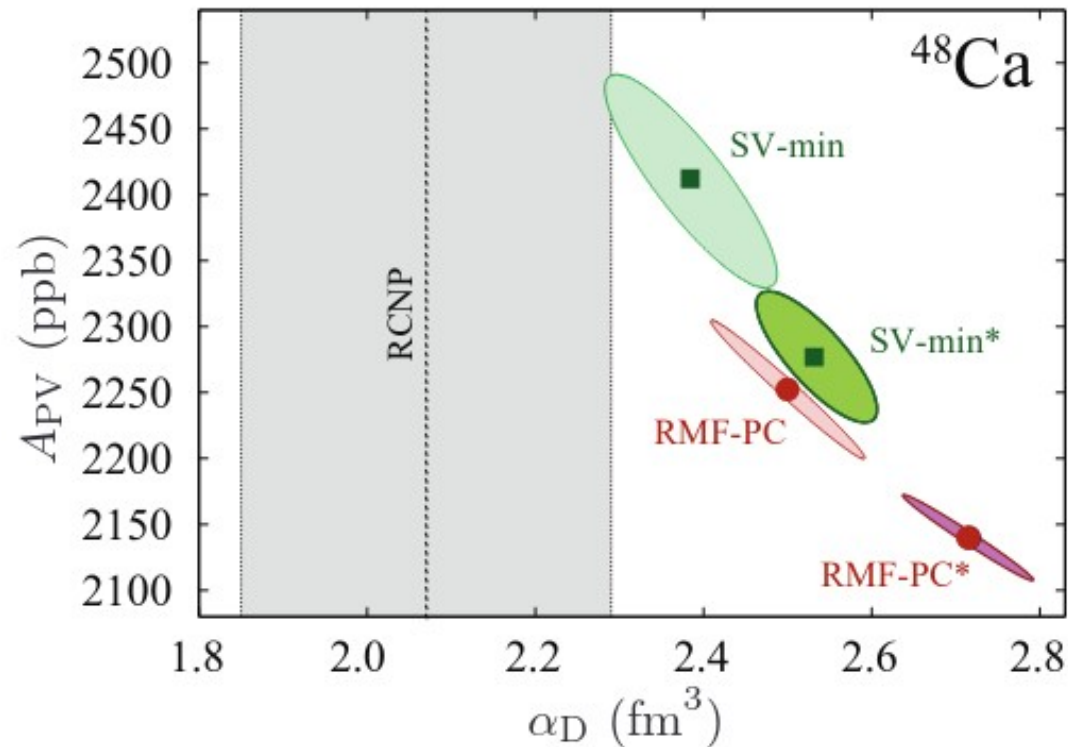
$$r_{\text{skin}} = 0.19 \pm 0.02 \text{ fm}$$

$$L = 54 \pm 8 \text{ MeV}$$

B and **Rch** in 208Pb not as good as other EDFs

Theoretical and experimental 1 σ errors overlap for SV-min

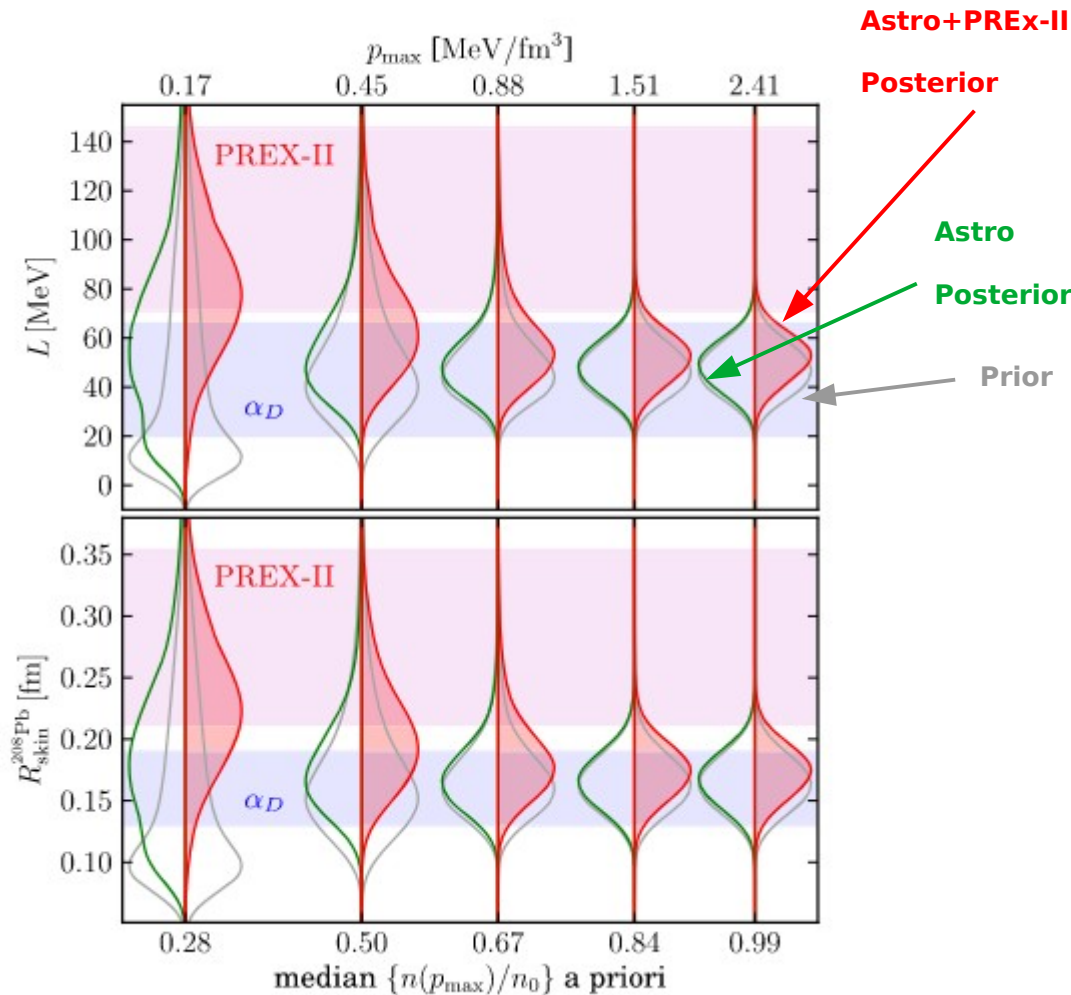
Predictions for CREx $A_{PV}(^{48}\text{Ca}) = 2400 \pm 60$ ppb.



Preliminary result from CREx $A_{pv} = 2568.6 \pm 113.2$ ppb

Our prediction is 1.5σ away from experiment

Other example: minimal model assumptions and statistical analysis



Non-parametric **equation of state** representation derived from **observations of neutron stars** with minimal modeling assumptions.

The resulting **astrophysical constraints** from heavy pulsar masses, LIGO/Virgo, and NICER clearly **favor “small”** values of the **neutron skin** and **L**.

Combining astrophysical data with **PREX-II** and **chiral effective field theory** constraints yields

$$J = 33.0 \pm 2.0 \text{ MeV}$$

$$L = 53 \pm 15 \text{ MeV}$$

$$R_{\text{skin}} = 0.17 \pm 0.04 \text{ fm}$$

Reed Essick, et al. Phys. Rev. Lett. 127, 192701 (2021)

Conclusions

- **Current EDFs** show strong *systematic* and *statistic* correlations between A_{pv} and ΔR_{np} or α_D and ΔR_{np}
- **Fitting masses and radii** do not give enough **information** on A_{pv} , α_D and ΔR_{np} (additional reason for the strong correlation)
- **Extending the fitting protocol** to include A_{pv} and α_D do **not change** the **correlations** above since **experimental errors** on these observables **are still large** → **model predictions** remain **biased** by **masses** and **radii**
- **Current EDFs** are able to overlap within $\sim 1\sigma$ all experimental **data** (preliminary: with the exception of A_{pv} in ^{48}Ca that is about 1.5σ away)
- More **accurate measurements** on A_{pv} could point to **model deficiencies**, while **EDFs** seem to **accommodate** better α_D (if also looking to other nuclei)
- **Ideally:** measure A_{pv} at **different kinematics** or A_{pv} on **different nuclei** with same (or better) accuracy → better **constraints** to **models**

Collaborators

B. K. Agrawal¹
P. F. Bortignon, M. Brenna,
G. Colò^{2,3}
W. Nazarewicz^{4,5,6}
N. Paar, D. Vretenar⁷
J. Piekarewicz⁸

P.-G. Reinhard⁹
Michal Warda¹⁰
Mario Centelles, Xavier
Viñas¹¹